

A new CP violating observable for the LHC

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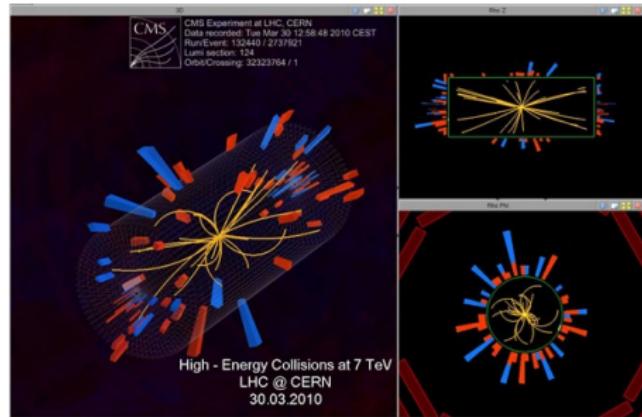
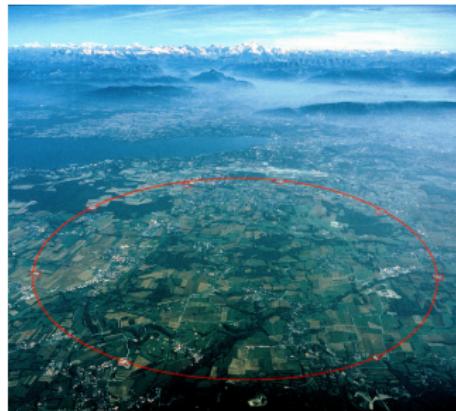
Cornell University

with Monika Blanke and Yuval Grossman

JHEP 1108 (2011) 033 (arXiv:1105.0672)

September 1, 2011

The LHC era has begun!



- ① Identify new states
- ② Measure masses and spins
- ③ Measure couplings, flavor structure, CP-violation

The Upshot

Goal: Find calculable & measurable \mathcal{CP} observables

- Requires interference & different strong phases
- So far: strong rescattering ($B \rightarrow K\pi$) and oscillation (meson mixing)
- Our result: new type of strong phase in 3-body decays with different orderings

Seeing CP-violation: Problem

- Looking for asymmetry:

$$\mathcal{A}_{\text{CP}} = \frac{\Gamma(i \rightarrow f) - \Gamma(\bar{i} \rightarrow \bar{f})}{\Gamma(i \rightarrow f) + \Gamma(\bar{i} \rightarrow \bar{f})} \neq 0$$

$$\mathcal{M} = |a| e^{i\varphi}$$

$$\overline{\mathcal{M}} = |a| e^{-i\varphi}$$

then

$$\mathcal{A}_{\text{CP}} = 0$$

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$$\mathcal{M} = |a_1|e^{i\varphi_1} + |a_2|e^{i\varphi_2}$$

$$\overline{\mathcal{M}} = |a_1|e^{-i\varphi_1} + |a_2|e^{-i\varphi_2}$$

then

$$\mathcal{A}_{\text{CP}} = 0$$

Seeing CP-violation: Problem

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$$\mathcal{A}_{\text{CP}} = \frac{\Gamma(i \rightarrow f) - \Gamma(\bar{i} \rightarrow \bar{f})}{\Gamma(i \rightarrow f) + \Gamma(\bar{i} \rightarrow \bar{f})} \neq 0$$

$$\mathcal{M} = |a_1| e^{i\delta_1 + i\varphi_1} + |a_2| e^{i\delta_2 + i\varphi_2}$$

$$\overline{\mathcal{M}} = |a_1| e^{i\delta_1 - i\varphi_1} + |a_2| e^{i\delta_2 - i\varphi_2}$$

then

$$\mathcal{A}_{\text{CP}} \propto |a_1||a_2|\sin(\delta_1 - \delta_2)\sin(\varphi_1 - \varphi_2)$$

Seeing CP-violation: Solution

- Requirements:

- ➊ Two interfering amplitudes a_1, a_2
- ➋ Different **weak** (CP-odd) phases φ_1, φ_2
- ➌ Different **strong** (CP-even) phases δ_1, δ_2

$$\mathcal{A}_{\text{CP}} \propto |a_1||a_2|\sin(\varphi_1 - \varphi_2)\sin(\delta_1 - \delta_2)$$

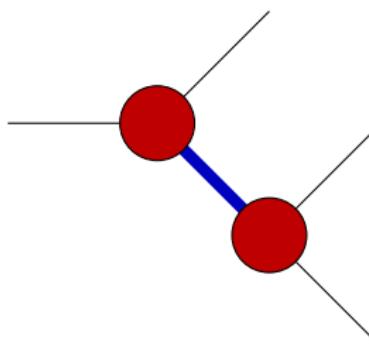
Strong phase?

- In general, comes from time evolution: e^{iEt}
- Basic case: oscillation of intermediate states - requires states with same quantum #'s
- More complicated: strong interaction rescattering - hard to calculate

Another way to get a **calculable strong phase?**

The Breit-Wigner Formula

- Process with narrow-width virtual state:


$$= \mathcal{M}_1 \frac{1}{q^2 - m^2 + i\Gamma m} \mathcal{M}_2$$

- Breit-Wigner propagator contributes phase
- Momentum-space equivalent of e^{iEt}

Strong phase from the propagator

Strong phase from intermediate particle:

- ① Different particles \leftrightarrow Time-integrated oscillation
- ② Different virtuality \rightarrow New!

$$\delta = \arg \left(\frac{1}{q^2 - m^2 + i m \Gamma} \right)$$

A new calculable strong phase

Requirements:

- ① Three body decay
- ② Two different orderings
- ③ On-shell resonance

Result:

CP-asymmetry in Dalitz plot

Toy model content

- All particles are scalars
- Heavy neutral particle: X_0^0
- Charged resonance: Y^+
- Lighter particles: $X_{1,2}^+, X_3^0$
- Phase space \implies scale hierarchy:



$$m_{X_0^0} > m_{Y^\pm} > m_{X_3^0} + m_{X_{1,2}^\pm}$$

Feynman rules

$$\begin{array}{ccc} & X_i^- & \\ & \swarrow \quad \searrow & \\ X_0^0 & \text{---} & X_3^0 \\ & \searrow \quad \swarrow & \\ & Y^+ & \end{array}$$
$$= -iae^{i\varphi_a} \qquad \qquad = -ibe^{i\varphi_b}$$

- One **weak** phase: $\varphi = \varphi_b - \varphi_a$

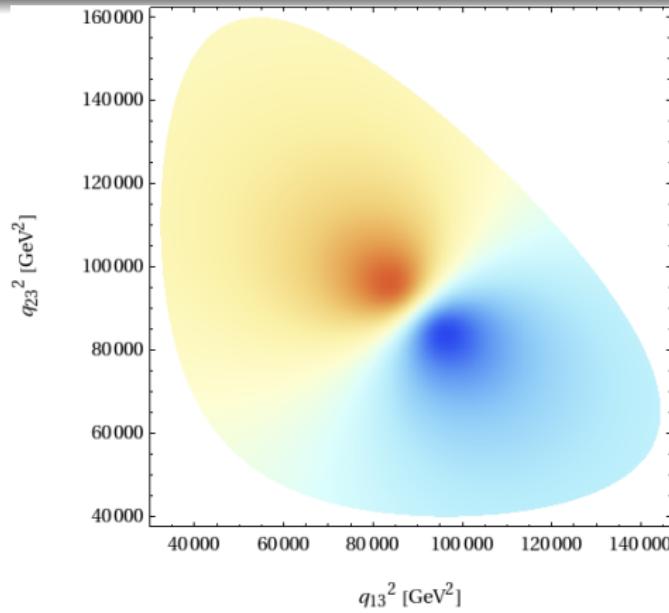
Toy model decays

$$X_0^0 \xrightarrow{\text{dashed}} X_1^+ \quad X_0^0 \xrightarrow{\text{dashed}} X_2^-$$
$$X_0^0 \xrightarrow{\text{dashed}} Y^- \quad X_0^0 \xrightarrow{\text{dashed}} Y^+$$
$$X_0^0 \xrightarrow{\text{dashed}} X_3^0 \quad X_0^0 \xrightarrow{\text{dashed}} X_3^0$$
$$= \frac{|a||b|e^{i\varphi}}{q_{23}^2 - m_Y^2 + im_Y\Gamma_Y}$$
$$= \frac{|a||b|e^{-i\varphi}}{q_{13}^2 - m_Y^2 + im_Y\Gamma_Y}$$

Different weak phase, different strong phase

Asymmetry in the Dalitz plot

$$\mathcal{A}_{\text{CP}}^{\text{diff}} \propto \sin 2\varphi (q_{13}^2 - q_{23}^2) \Gamma_Y m_Y$$



Integrated asymmetries

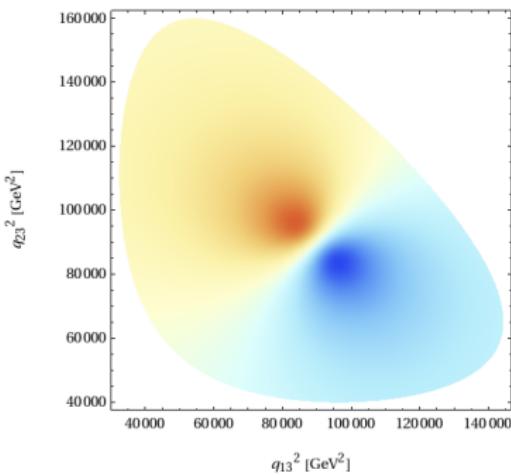
$$X_0^0 \rightarrow X_1^+ X_2^- X_3^0$$

- Integrated rate suppressed:

$$\mathcal{A}_{\text{CP}}^{\text{int}} \propto \frac{\Delta m_{12}^2}{m_0^2}$$

- Eliminate suppression by phase space weighting:

$$\mathcal{A}_{\text{CP}}^{\text{wgt}} \equiv \frac{1}{\Gamma + \bar{\Gamma}} \int dq_{13}^2 dq_{23}^2 \operatorname{sgn}(q_{23}^2 - q_{13}^2) \left(\frac{d\Gamma}{dq_{13}^2 dq_{23}^2} - \frac{d\bar{\Gamma}}{dq_{13}^2 dq_{23}^2} \right)$$



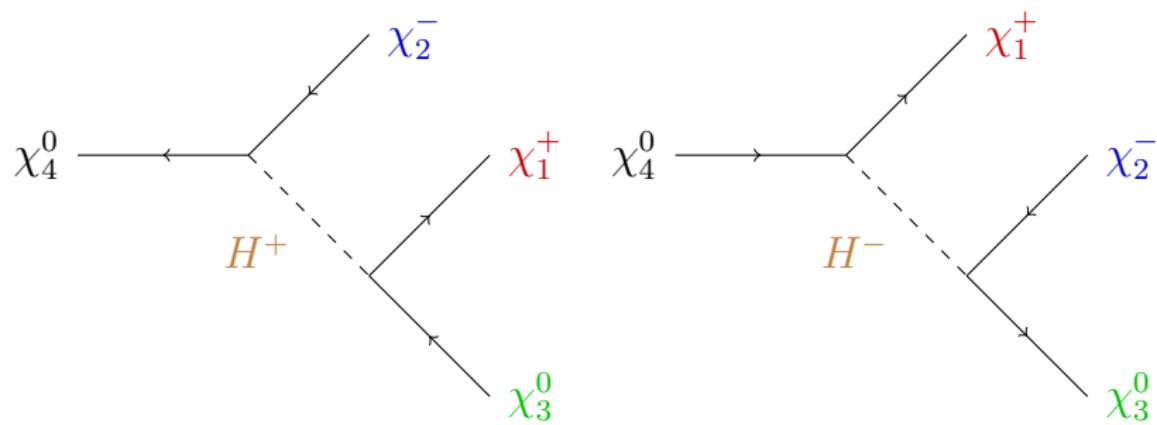
The relevant model

Electroweak sector of MSSM

- Heavy neutral particle: $\sim \tilde{B}$
- Intermediate charged resonance: H^\pm
- “Light” final states: lighter charginos and neutralinos
- Hierarchy of scales for maximal signal:

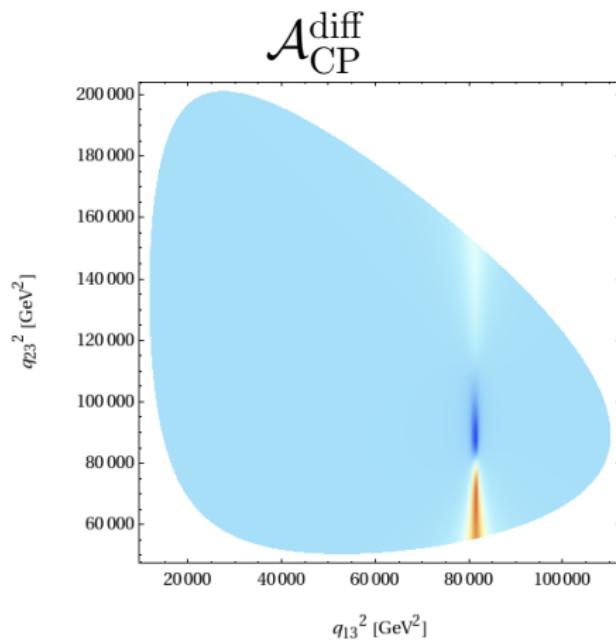
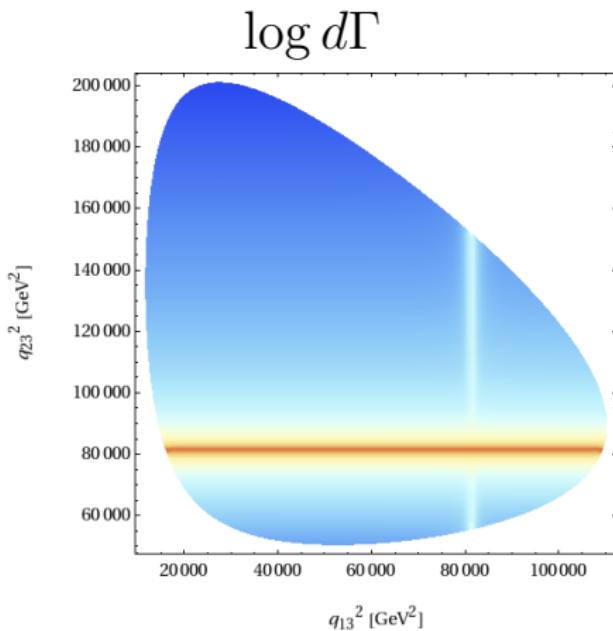
$$m_{\chi_4^0} \sim M_1 \gg m_{H^\pm} \gg m_{\chi_i^0}, m_{\chi_j^\pm} \sim \sqrt{|\mu M_2|} > m_Z$$

The Feynman diagrams



- One weak phase: $\arg(\mu b^* M_2)$

Dalitz plot observables



MSSM results

- Suppressed integrated asymmetry:

$$\mathcal{A}_{\text{CP}}^{\text{int}} = -3.5 \times 10^{-5}$$

- Using phase space weighting:

$$\mathcal{A}_{\text{CP}}^{\text{wgt}} = -6.5 \times 10^{-4}$$

Electroweak MSSM is challenging

The ingredients

Recipe for Dalitz plot asymmetry:

- Three body decay
- Two different orderings
- On-shell resonance

